# The Theoretical Minimum 

Classical Mechanics - Solutions
L05E03
Last version: tales.mbivert.com/on-the-theoretical-minimum-solutions/ or github.com/mbivert/ttm
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May 10, 2023

Exercise 1. Rework Exercise 2 for the potential $V=\frac{k}{2\left(x^{2}+y^{2}\right)}$. Are there circular orbits? If so, do they all have the same period? Is the total energy conserved?

## Equations of motion

The approach is similar to what has been done for the previous exercise: for this system, the potential energy $V$ is:

$$
\begin{equation*}
V=\frac{k}{2\left(x^{2}+y^{2}\right)} \tag{1}
\end{equation*}
$$

By Newton's second law of motion ${ }^{1}$, given $\boldsymbol{r}=(x, y)$, we have:

$$
\begin{equation*}
\boldsymbol{F}=m \boldsymbol{a}=m \dot{\boldsymbol{v}}=m \ddot{\boldsymbol{r}} \tag{2}
\end{equation*}
$$

Or,

$$
\begin{align*}
& F_{x}=m \ddot{x}  \tag{3}\\
& F_{y}=m \ddot{y}
\end{align*}
$$

We know by equation (5) of this lecture that to each coordinate $x_{i}$ of the configuration space $\{x\}$, there is a force $F_{i}$, derived from the potential energy $V$ :

$$
\begin{equation*}
F_{i}(\{x\})=-\frac{\partial}{\partial x_{i}} V(\{x\}) \tag{4}
\end{equation*}
$$

As for the previous exercise, we make heavy use of the chain rul ${ }^{2}$ for derivation:

$$
\begin{equation*}
\frac{d}{d x} f(g(x))=g^{\prime}(x) f^{\prime}(g(x)) \tag{5}
\end{equation*}
$$

To compute e.g. $F_{x}(x, y)$, we define $\phi(x)=x^{2}+y^{2}$ :

$$
\begin{align*}
F_{x}(x, y) & =-\frac{\partial}{\partial x} V(x, y) \\
& =\frac{k}{2} \frac{d}{d x} \frac{1}{\phi(x)} \\
& =\frac{k}{2} \phi^{\prime}(x) \frac{-1}{\sqrt{\phi(x)}}  \tag{6}\\
& =\frac{k x}{\left(x^{2}+y^{2}\right)^{2}}
\end{align*}
$$

Thus finally:

[^0]\[

$$
\begin{align*}
& F_{x}(x, y)=\frac{k x}{\left(x^{2}+y^{2}\right)^{2}} \\
& F_{y}(x, y)=\frac{k y}{\left(x^{2}+y^{2}\right)^{2}} \tag{7}
\end{align*}
$$
\]

Hence combining (7) and (3):

$$
\begin{align*}
& F_{x}(x, y)=m \ddot{x}(t)=k \frac{x(t)}{\left(x(t)^{2}+y(t)^{2}\right)^{2}} \\
& F_{y}(x, y)=m \ddot{y}(t)=k \frac{y(t)}{\left(x(t)^{2}+y(t)^{2}\right)^{2}} \tag{8}
\end{align*}
$$

## Circular orbits

Let's make a guess, and see what would happen were we to plug the simplest circular motion, that we've already studied in the book at the end of Chapter 2 (Motion), given by:

$$
x(t)=R \cos (\omega t) ; \quad y(t)=R \sin (\omega t)
$$

Which is very convenient for us, because if we try this solution in 88, the (common) denominator simplifies:

$$
\left(x(t)^{2}+y(t)^{2}\right)^{2}=\left((R \cos (\omega t))^{2}+(R \sin (\omega t))^{2}\right)^{2}=R^{4} \underbrace{\left(\cos ^{2}(\omega t)+\sin ^{2}(\omega t)\right)^{2}}_{=1}=R^{4}
$$

Let's now consider the velocities and accelerations we would obtain by differentiating our guess for $x(t)$ and $y(t)$ :

$$
\begin{array}{rlr}
\dot{x}(t)= & -R \omega \sin (\omega t) ; & \dot{y}(t)= \\
\ddot{x}(t)= & -R \omega^{2} \cos (\omega t) ; & \ddot{y}(t)= \\
& -R \omega^{2} \sin (\omega t) \\
\end{array}
$$

There are two ways for this guess to actually work:

1. Either we set $\omega^{2}=-k / m R^{4}$, which implies either:

- $k$ to be zero (trivial solution then);
- or that $m R$ to be close to infinite (unrealistic);
- or that $k$ is (strictly) negative;
- or that either $m$ or $R$ are negative (unrealistic);
- or, mathematically, that $\omega$ is an imaginary (complex) number, which would be difficult to interpret, physically;

2. The other option would be for $R$ to be negative, which again doesn't make a lot of sense, physicallywise.

Remark 1. Note that our guess would have worked for a negated $V$ :

$$
V=-\frac{k}{2\left(x^{2}+y^{2}\right)}
$$

Remark 2. What is commonly referred to as "the trivial solution", especially in the context of differential equations, is the solution $x(t)=0$, which is of little interest, mathematically and physically.

We can conclude that, at least physically speaking, there are no circular orbits, unless $k$ is negative. This is because, if there were circular orbits, then they would be a coordinate change away from being in the form of our guess.

The only remaining issue is that $k$ hasn't been clearly defined, physically speaking, so we can't really know for sure if assuming $k$ to be negative (with a reminder that $k=0$ leads to the trivial solution).

Remark 3. Another approach, used for instance in the official solution $4^{3}$, relies on the polar coordinate $(r, \theta)$ : the existence of a circular orbit then translate to $r$ being a constant, or equivalently, $\dot{r}=0$.

We'll dive deeper into polar coordinates in a later exercise, alongside a bunch of other elements related to circular motion ( L06E05, which involves a pendulum).

## Energy conservation

Earlier in the lecture, the kinetic energy has been defined to be the sum of all the kinetic energies for each coordinate:

$$
\begin{equation*}
T=\frac{1}{2} \sum_{i} m_{i}{\dot{x_{i}}}^{2} \tag{9}
\end{equation*}
$$

Which gives us for this system, expliciting the time-dependencies:

$$
\begin{equation*}
T(t)=\frac{1}{2} m \dot{x}(t)^{2}+\frac{1}{2} m \dot{y}(t)^{2}=\frac{1}{2} m\left(\dot{x}(t)^{2}+\dot{y}(t)^{2}\right) \tag{10}
\end{equation*}
$$

From which we can compute the variation of kinetic energy over time, again using the chain rule:

$$
\begin{align*}
\frac{d}{d t} T(t) & =\frac{1}{2} m(2 \dot{x}(t) \ddot{x}(t)+2 \dot{y}(t) \ddot{y}(t))  \tag{11}\\
& =m(\dot{x} \ddot{x}+\dot{y} \ddot{y})
\end{align*}
$$

On the other hand, we can compute the variation of potential energy over time from (1). We'll use the chain rule again, with $\phi(t)=x(t)^{2}+y(t)^{2}$ and thus:

$$
\begin{aligned}
\phi^{\prime}(t) & =2 x^{\prime}(t) x(t)+2 y^{\prime}(t) y(t) \\
& =2 \dot{x} x+2 \dot{y} y
\end{aligned}
$$

It follows that:

$$
\begin{align*}
\frac{d}{d t} V(t) & =\frac{d}{d t} \frac{k}{2\left(x(t)^{2}+y(t)^{2}\right)} \\
& =\frac{k}{2} \frac{d}{d t} \phi(t)^{-1} \\
& =-\frac{k}{2} \phi^{\prime}(t) \phi(t)^{-2} \\
& =-\frac{k}{2} \frac{2 \dot{x} x+2 \dot{y} y}{\left(x(t)^{2}+y(t)^{2}\right.}  \tag{12}\\
& =-k \frac{\dot{x} x+\dot{y} y}{\left(x^{2}+y^{2}\right)^{2}} \\
& =-k \frac{\dot{x} x+\dot{y} y}{\phi(t)^{2}}
\end{align*}
$$

Then, from (8), we can extract

$$
x(t)=\frac{m}{k} \ddot{x} \phi(t)^{2} ; \quad y(t)=\frac{m}{k} \ddot{y} \phi(t)^{2}
$$

Injecting in (12) gives:

$$
\begin{align*}
\frac{d}{d t} V(t) & =-\frac{k}{\phi(t)^{2}}\left(\dot{x} \frac{m}{k} \ddot{x} \phi(t)^{2}+\dot{y} \frac{m}{k} \ddot{y} \phi(t)^{2}\right)  \tag{13}\\
& =-m(\dot{x} \ddot{x}+\dot{y} \ddot{y})
\end{align*}
$$

And so by combining (13) and we can indeed see that the energy is conserved:

$$
\frac{d}{d t} E(t)=\frac{d}{d t} T(t)+\frac{d}{d t} V(t)=0
$$

[^1]
[^0]:    1 https://en.wikipedia.org/wiki/Newton\%27s_laws_of_motion\#Second
    2 https://en.wikipedia.org/wiki/Chain_rule

[^1]:    $\sqrt[3]{\text { http://www.madscitech.org/tm/slns/l5e3.pdf }}$

