

The Theoretical Minimum

Classical Mechanics - Solutions

L11E04

Last version: tales.mbvivert.com/on-the-theoretical-minimum-solutions/ or github.com/mbvivert/ttm

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Exercise 1. Using the Hamiltonian, Eq. (24), work out Hamilton's equations of motion and show that you just get back to the Newton-Lorentz equation of motion.

The Hamiltonian is:

$$H = \frac{1}{2m} \sum_i \left(p_i - \frac{e}{c} A_i(\mathbf{q}) \right)^2$$

Hamilton's equation of motion are given by the pair:

$$\dot{q}_j = \frac{\partial H}{\partial p_j}; \quad \dot{p}_j = -\frac{\partial H}{\partial q_j}$$

Let's get started with the first one:

$$\begin{aligned} \dot{q}_j &= \frac{\partial H}{\partial p_j} \\ &= \frac{1}{2m} \frac{\partial}{\partial p_j} \sum_i \left(p_i - \frac{e}{c} A_i(\mathbf{q}) \right)^2 \\ &= \frac{1}{2m} \frac{\partial}{\partial p_j} \sum_i \left(p_i^2 - 2\frac{e}{c} p_i A_i(\mathbf{q}) + \left(\frac{e}{c}\right)^2 A_i(\mathbf{q})^2 \right) \\ &= \frac{1}{2m} \sum_i \left(\frac{\partial p_i^2}{\partial p_j} - 2\frac{e}{c} \left(\frac{\partial p_i}{\partial p_j} A_i(\mathbf{q}) + p_i \underbrace{\frac{\partial A_i(\mathbf{q})}{\partial p_j}}_{=0} \right) + \left(\frac{e}{c}\right)^2 \underbrace{\frac{\partial A_i(\mathbf{q})^2}{\partial p_j}}_{=0} \right) \\ &= \frac{1}{2m} \sum_i \left(\frac{\partial p_i^2}{\partial p_j} - 2\frac{e}{c} \frac{\partial p_i}{\partial p_j} A_i(\mathbf{q}) \right) \\ &= \frac{1}{2m} \left(2p_j - 2\frac{e}{c} A_j(\mathbf{q}) \right) \\ &= \boxed{\frac{1}{m} (p_j - \frac{e}{c} A_j(\mathbf{q}))} \end{aligned}$$

We've found back the expression of the moment, tweaked a little. Let's now take the time derivative of the previous equation to get:

$$\begin{aligned} \ddot{q}_j &= \frac{1}{m} \left(\dot{p}_j - \frac{e}{c} \dot{A}_j(\mathbf{q}) \right) \\ \Leftrightarrow m\ddot{q}_j &= \dot{p}_j - \frac{e}{c} \dot{A}_j(\mathbf{q}) \end{aligned}$$

Which starts to look like an equation of motion. We'll develop $\dot{A}_j(\mathbf{q})$ using the chain rule later. For now, let's compute \dot{p}_j , but this time, let's make our life a bit simpler by using the chain rule:

$$\begin{aligned}
\dot{p}_j &= -\frac{\partial H}{\partial q_j} \\
&= -\frac{1}{2m} \frac{\partial}{\partial q_j} \sum_i \left(p_i - \frac{e}{c} A_i(\mathbf{q}) \right)^2 \\
&= -\frac{1}{2m} \sum_i 2\phi \frac{\partial \phi}{\partial q_j} \\
&= \sum_i \underbrace{\frac{1}{m} \left(p_i - \frac{e}{c} A_i(\mathbf{q}) \right)}_{=q_i} \frac{e}{c} \frac{\partial A_i(\mathbf{q})}{\partial q_i} \\
&= \boxed{\frac{e}{c} \sum_i \dot{q}_i \frac{\partial A_i(\mathbf{q})}{\partial q_i}}
\end{aligned}$$

Now let's use the multi-dimensional chain rule to develop $\dot{A}_j(\mathbf{q})$:

$$\dot{A}_j(\mathbf{q}) = \dot{A}_j(q_x(t), q_y(t), q_z(t)) = \sum_i \frac{\partial A_j(\mathbf{q})}{\partial q_i} \dot{q}_i$$

Finally, let's use this and the expression of p_j to rewrite our embryo of motion equation:

$$\begin{aligned}
m\ddot{q}_j &= \dot{p}_j - \frac{e}{c} \dot{A}_j(\mathbf{q}) \\
&= \sum_i \dot{q}_i \frac{e}{c} \frac{\partial A_i(\mathbf{q})}{\partial q_i} - \frac{e}{c} \sum_i \frac{\partial A_j(\mathbf{q})}{\partial q_i} \dot{q}_i \\
&= \frac{e}{c} \sum_i \dot{q}_i \left(\frac{\partial A_i(\mathbf{q})}{\partial q_i} - \frac{\partial A_j(\mathbf{q})}{\partial q_i} \right) \\
&= \frac{e}{c} \left(\underbrace{\dot{q}_j \left(\frac{\partial A_j(\mathbf{q})}{\partial q_j} - \frac{\partial A_j(\mathbf{q})}{\partial q_j} \right)}_{=0} + \dot{q}_k \underbrace{\left(\frac{\partial A_k(\mathbf{q})}{\partial q_k} - \frac{\partial A_j(\mathbf{q})}{\partial q_k} \right)}_{=B_l} + \dot{q}_l \underbrace{\left(\frac{\partial A_l(\mathbf{q})}{\partial q_l} - \frac{\partial A_j(\mathbf{q})}{\partial q_l} \right)}_{=-B_k} \right) \\
&= \boxed{\frac{e}{c} (B_l \dot{q}_k - B_k \dot{q}_l)}
\end{aligned}$$

And so for j, k, l three distinct elements of $\{x, y, z\}$, we indeed have found our equations of Newton-Lorentz:

$$\boxed{m\ddot{q}_j = \frac{e}{c} (B_l \dot{q}_k - B_k \dot{q}_l)}$$